



Full length article



## The origin of phase separation in binary aluminosilicate glasses

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### ABSTRACT

The quest for hard and tough transparent oxide glasses is at the core of glass science and technology. Aluminosilicate glasses exhibiting nanoscale phase separation emerge as promising candidates for such materials. Nevertheless, proper control of the phase separation represents a daunting challenge due to its elusive origins. Here we employ large-scale molecular dynamics simulations and structural analysis to unravel the underlying mechanisms of the phase separation in aluminosilicate. The observed phase separation originates from an arrangement of  $\text{SiO}_4$  and  $\text{AlO}_n$  polyhedra, which manifests from the second coordination shell and extends to higher shells. This specific arrangement is driven by repulsion between the polyhedra, reaching its maximum at around 50 mol% of  $\text{Al}_2\text{O}_3$ . This behavior becomes pronounced around and below the glass transition temperature. This work sheds light on the origin of phase separation and provides a route for further exploration across other compositions to develop glasses with adapted mechanical performance.

Oxide glasses are used in a wide range of applications, enabling modern civilization, e.g., optical fibers [1], bioactive materials [2,3], solar cells [4], and flexible and foldable screens for mobile devices [5]. The usage of oxide glasses in these applications benefits from their excellent optical properties [6,7], chemical durability [8], and the high tailorability of the properties through compositional design [2,9], (post-) processing [10,11], or a combination thereof [12]. However, the brittleness of these glasses, which is due to the lack of a stable shearing mechanism that accommodates the deformation, has limited the development of existing and emerging glass applications [13]. Thus, the need for new glasses with adapted mechanical properties triggered efforts to achieve this goal [14–19].

For instance, the use of reinforcements has been proven effective in controlling crack propagation [20], but it affects the glass transparency [21]. Glasses with nanoscale phase separation were observed to improve fracture toughness compared to homogeneous glasses [22, 23]. The mechanism of enhancing the hardness and fracture toughness of glasses with phase separation works in principle in the same way as for a composite material with two interconnected glassy regions with different structures, i.e., a rigid one and the other one is more flexible [18,23,24].

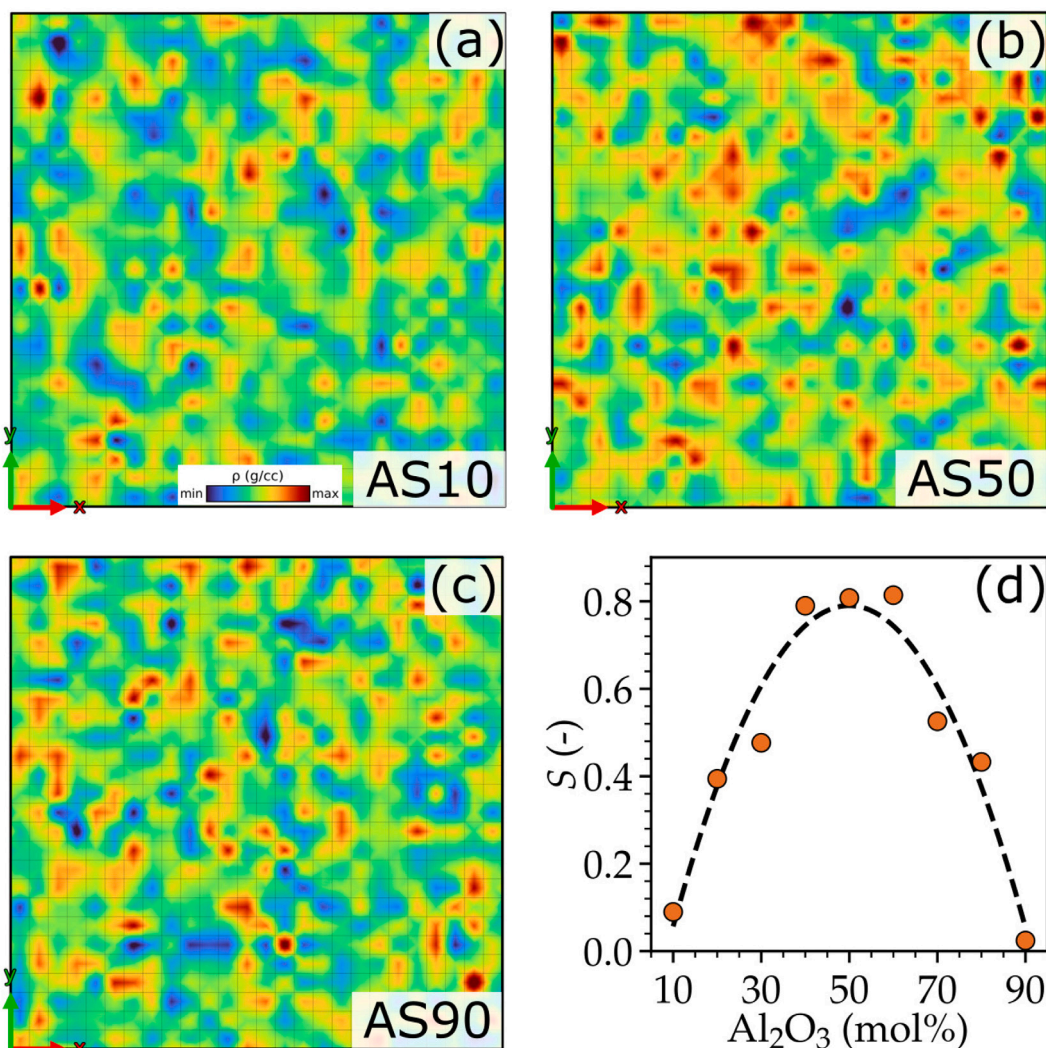
Binary aluminosilicate glasses with nanoscale phase separation are great potential candidates to make glasses with high hardness and crack resistance [19,25]. Rosales-Sosa et al. studied aluminosilicate glasses with alumina content ranging between 30 and 60 mol% and showed that the increase in the elastic moduli, hardness, and cracking

resistance of binary aluminosilicate glasses was due to the increase in the atomic packing density, the high dissociation energy per unit volume of  $\text{Al}_2\text{O}_3$  compared to that of  $\text{SiO}_2$ , and increase in the glass density with increasing  $\text{Al}_2\text{O}_3$  content [19]. While the explanations from Ref. [19] are indeed valid, they neglect the presence of a phase separation, which shows a maximum at the compositions with an alumina content around 50 mol% [26–32], and its effects on the observed behavior. Macdowell et al. highlighted that the composition at which the alumina content is around 50 mol% shows the highest degree of phase separation [26]. While others showed that the maximum phase separation occurs around 60 mol% of  $\text{Al}_2\text{O}_3$  [31,33]. Recently, Urata et al. [34] used molecular dynamics simulations to study the co-doping effect of alumina and fluorine on the density fluctuations in silica glass. In their study, they also observed phase separation in binary aluminosilicate glasses in the simulated samples. Although they reported phase separation, no indication of its origins or deeper insights was given.

Although phase separation has been observed in binary [34–37], ternary [38] aluminosilicate, and other oxide glasses [39], its origins remain elusive. Here we used a million atoms molecular dynamics (MD) simulations to study the phase separation in aluminosilicate glasses with  $\text{Al}_2\text{O}_3$  content ranging between 10 and 90 mol%. We found that this phase separation is discernible through changes in the medium-range structure, starting from the second coordination shell

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**Fig. 1.** (a–c) Selected local density maps of binary aluminosilicate glasses. The density maps are for (a)  $x = 10$  mol%, (b)  $x = 50$  mol%, and (c)  $x = 90$  mol%. The simulation boxes were divided to produce small cubes of volume,  $10 \times 10 \times 10 \text{ \AA}^3$ , and then the density of each cube was calculated and used for further analysis. In (d), the skewness is plotted as a function of the composition at 300 K. The skewness is obtained by fitting the histograms of local density to a skewed Gaussian function (See Supplementary material for details). The dashed line is fit to a quadratic function and serves as a guide to the eye (For interpretation of the colors in the figure(s), the reader is referred to the web version of this article).

and extending to larger shells. Furthermore, our MD results indicate an absence of any significant phase separation in the liquid state at elevated temperatures, and signs of phase separation were observed only in the supercooled liquid around the glass transition temperature ( $T_g$ ) and below it.

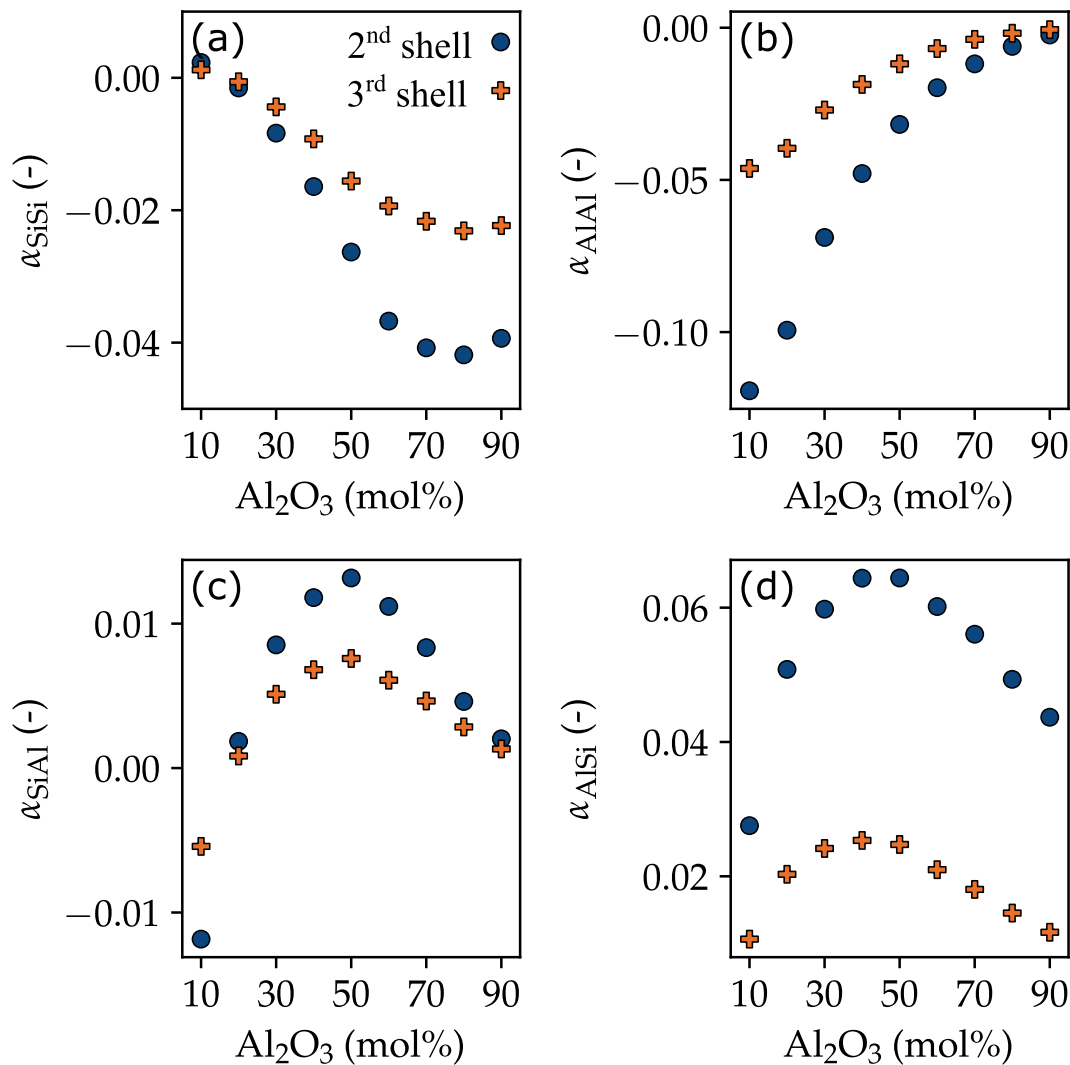
All MD simulations were performed using LAMMPS [40]. The potential from Bouhadja et al. [41,42] was used to model the interactions between atoms. The glass composition of the binary aluminosilicate glasses has  $\text{Al}_2\text{O}_3$  content ranging between 10 mol% to 90 mol%, with samples named  $\text{AS}_x$  with  $x$  is the molar content of  $\text{Al}_2\text{O}_3$ . The glass was equilibrated at 5000 K and cooled using 1 K/ps to room temperature (300 K). Configurations at different temperatures during the cooling process were extracted for further analysis. More details about the glass preparation can be found in Ref. [43] and the supplementary materials (SM). The visualizations are made using OVITO [44]. The structure of the glass is in realistic agreement with the data available in the literature [41,45,46] as discussed in the SM and shown in Fig. S2, Fig. S3, Fig. S4, and Tab. S1.

The phase separation was visually observed in the simulated aluminosilicate glasses with different degrees as illustrated from selected compositions in Fig. 1(a–c) using the 2D spatial distribution of the mass density. The histograms of the density distribution throughout

the system are shown in the SM, Fig. S5, which were fitted using a skewed Gaussian distribution function (details in the SM) to extract the skewness  $S$ . If the glass is homogeneous, then the density distribution should be well described through a Gaussian function, with its skewness being nearly 0. On the other hand, if the glass has phase separation, the density distribution can no longer be described through a symmetric Gaussian function, and it will be skewed. The higher the values of the skewness, the higher the asymmetry of the distribution. The values of the skewness  $S$  are plotted as a function of the alumina content and show a maximum at composition around 50 mol% of  $\text{Al}_2\text{O}_3$ .

The glasses are expected to be more homogeneous at the compositions at which the  $\text{SiO}_2$  content is very high or very low. The presence of inverted parabolic-like behavior of  $S$  with the alumina content and the occurrence of a maximum at around 50 mol% of  $\text{Al}_2\text{O}_3$  is consistent with the hypothesis stating that the maximum phase separation is around 50 mol% of  $\text{Al}_2\text{O}_3$ . These results are in good agreement with experiments and theoretical studies, although there is disagreement regarding the value of the critical composition of demixing in binary AS glasses [27,29,47].

From an atomistic point of view, the tendency to phase separate in AS glasses could be correlated with the preference of  $\text{SiO}_4$  and  $\text{AlO}_n$  (with  $n = 4, 5$ , or 6) polyhedra to be surrounded by those of the same



**Fig. 2.** Chemical short-range order at the second and third shell as a function of the alumina content at 300 K. (a) Si-Si, (b) Al-Al, (c) Al-Si, and (d) Si-Al. The standard error bars are smaller than the symbol size. We note that an exact determination of the higher coordination shells in amorphous materials is far from trivial. Thus,  $r_{cut}^{2nd\ shell} = 5.35 \text{ \AA}$  and  $r_{cut}^{3rd\ shell} = 9.1 \text{ \AA}$ . The values of the cutoff radii are provided for better reproducibility of the results. (For interpretation of the colors in the figure(s), the reader is referred to the web version of this article).

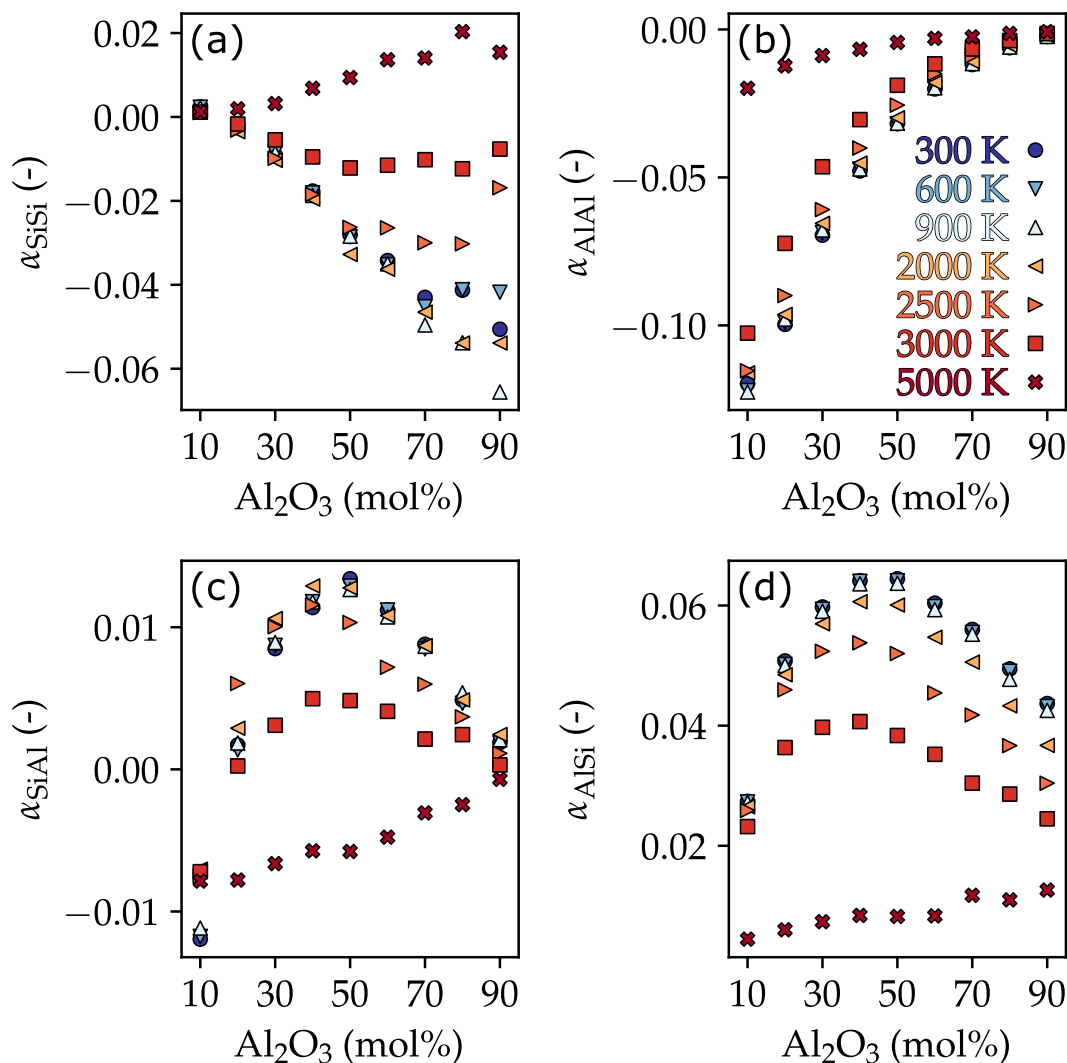
type. This process is evaluated with the distance-dependent Warren-Cowley (WC) order parameter defined for a pair of atom types  $i, j$  as follows [11,48]:

$$\alpha_{ij}(r) = 1 - \frac{n_{ij}(r)}{N_i(r)c_j}, \quad (1)$$

where  $n_{ij}(r)$  is the number of atoms of type  $j$  neighbors of a central atom of type  $i$  at a distance  $r$ ,  $N_i(r)$  is the total of neighbors for the atom of type  $i$  at a distance  $r$ ,  $c_j$  is the concentration of type  $j$ . Negative values of  $\alpha_{ij}$  indicate attraction between the pairs, while positive values correspond to the opposite. Throughout this paper, we limited the calculation of  $\alpha_{ij}$  to Si and Al atoms without involving oxygen atoms to extract the correlations between  $\text{SiO}_4$  and  $\text{AlO}_n$  polyhedra with  $n = 4, 5, \text{ or } 6$ . Fig. S6 represents the variation of this order parameter in the first coordination shell, defined by the first minimum of the radial distribution function, as a function of  $\text{Al}_2\text{O}_3$  content. The calculated  $\alpha_{\text{SiSi}}$  is slightly higher than zero up to  $x = 0.6$ , where it starts to increase. At low  $\text{Al}_2\text{O}_3$  content,  $\alpha_{\text{SiAl}}$  is negative and increases with increasing  $\text{Al}_2\text{O}_3$  content. This indicates a preference for Si to have Al in their first neighboring shell, which decreases with increasing alumina. On the other hand,  $\alpha_{\text{AlSi}}$  is always positive, highlighting that  $\text{AlO}_n$  polyhedra do not prefer  $\text{SiO}_4$  tetrahedra as neighbors and indicating a relatively low probability of having  $\text{SiO}_4$  tetrahedra around an  $\text{AlO}_n$  polyhedra in

the first shell. The WC parameter for Al-Al is negative at low alumina content and increases with increasing  $\text{Al}_2\text{O}_3$ . This indicates that there is a strong attraction between  $\text{AlO}_n$  and  $\text{AlO}_n$  polyhedra, and  $\text{AlO}_n$  prefers to have other  $\text{AlO}_n$  polyhedra in their first neighboring shell, which becomes less pronounced with increasing  $\text{Al}_2\text{O}_3$ .

Since the data obtained from the  $\alpha_{ij}$  up to the first coordination shell, involving neighboring polyhedra, do not show clear signs of phase separation, which is expected as the phase separation spans over larger length scales. The investigation of the chemical arrangement of the glass building blocks at larger coordination shells is necessary. In Fig. 2, the evolution of the WC parameters in the second and third coordination shells of all considered compositions is represented. Comparing the results obtained from Fig. S6, which focuses on the first shell exclusively, with those presented in Fig. 2, a significant change in the WC parameter is observed for the Si-Si pair.  $\alpha_{\text{SiSi}}$  calculated from the second shell is positive at the composition with 90 mol% of  $\text{SiO}_2$  and becomes negative with increasing  $\text{Al}_2\text{O}_3$  content. This indicates that the second and third shells of Si prefer to have Si than Al, which is the first indication of the phase separation. The  $\alpha_{\text{AlAl}}$  in the second and third shells remains negative and follows the same behavior as discussed earlier for the first shell. This finding indicates the preferred ordering between the  $\text{AlO}_n$  polyhedra. More interestingly,



**Fig. 3.** Change of  $\alpha_{ij}$  at the second shell as a function of the composition for different selected temperatures above and below  $T_g$ . (a) is for  $\alpha_{\text{SiSi}}$ , (b) is for  $\alpha_{\text{AlAl}}$ , (c) is for  $\alpha_{\text{SiAl}}$ , and (d) is for  $\alpha_{\text{AlSi}}$ . The data for all samples as a function of the temperature is shown in Fig. S7. The standard error bars are smaller than the symbol size. The  $r_{\text{cut}}^{\text{2nd shell}} = 5.35 \text{ \AA}$ . The values of the cutoff radii are provided for better reproducibility of the results. (For interpretation of the colors in the figure(s), the reader is referred to the web version of this article).

the WC parameter for the pairs Si–Al and Al–Si shows a maximum around  $x = 0.5$ , consistent with the data provided in Fig. 1(d). A gradual increase is observed for both pairs up to a maximum and decreases again. This shows that a maximum in the repulsion between  $\text{SiO}_4$  tetrahedra and  $\text{AlO}_n$  polyhedra is around the composition at which  $R = \text{SiO}_2/\text{Al}_2\text{O}_3 = 1$ .

The temperature-dependent evolution of the WC parameter  $\alpha_{\text{SiSi}}$ ,  $\alpha_{\text{AlAl}}$ ,  $\alpha_{\text{SiAl}}$ , and  $\alpha_{\text{AlSi}}$  for the second shell is shown in Fig. 3 and used to investigate at which temperature the phase separation is manifested in the samples across all compositions. At elevated temperatures, the behavior of  $\alpha_{\text{SiSi}}$  and  $\alpha_{\text{AlAl}}$  of the second neighboring shell is similar to that at 300 K, with only the intensity being different. On the other hand, and at high temperatures, it is evident from the graph that the WC parameter of the  $\alpha_{\text{SiAl}}$ , and  $\alpha_{\text{AlSi}}$  tends towards lower values, negative in case of  $\alpha_{\text{SiAl}}$ , and positive ones in case of  $\alpha_{\text{AlSi}}$ , which indicate that in the liquid phase,  $\text{SiO}_4$  attract  $\text{AlO}_n$ . The WC parameter becomes increasingly repulsive as the temperature decreases, particularly around the glass transition temperature (See Tab. S1 and Fig. S8). This suggests that aluminum polyhedra prefer to be surrounded by other aluminum polyhedra. This tendency promotes immiscibility in the supercooled liquid phase, leading to phase separation. The WC parameter shows that the sample exhibits a strong tendency towards demixing in the low-temperature region, as indicated by the higher positive values.

The findings presented in this study shed light on the origins of the nanoscale phase separation in binary aluminosilicate glasses. The repulsion between  $\text{SiO}_4$  and  $\text{AlO}_n$  polyhedra is shown to be the origin of the phase separation in aluminosilicate glasses at the atomic scale. This phase separation was observed in previous experimental studies [26–28,31,33] and the MD work of Urata et al. [34]. Our analysis showed that the composition at which the phase separation is maximized in binary AS glasses is around 50 mol% of  $\text{Al}_2\text{O}_3$ .

The underlying hypothesis of this discussion is that phase separation refers to the redistribution of atoms from a homogeneous mixture to form distinct regions with different local atomic arrangements. This process occurs to minimize the energy of the system and to accommodate the differences in the local potential energy between different regions of the system. We computed the atomic potential energy for Si and Al to verify this hypothesis. The atomic potential energies ( $\langle U_i \rangle$ ) were calculated for both Al and Si atoms. Subsequently, the variance of the potential energy per atom for each type of atom was calculated using the formula:  $\sigma_i^2 = (1/N_i)(U_i - \langle U_i \rangle)^2$ , where  $N_i$  represents the number of atoms of type  $i$  and  $U_i$  is the potential energy of atoms of type  $i$ . In Fig. 4(a), the variance obtained from the atomic potential energy distribution for both Si and Al atoms is shown. This shows that as the concentration of  $\text{Al}_2\text{O}_3$  increases, there is a corresponding non-linear decrease in the variance for Al with a change of slope around 50

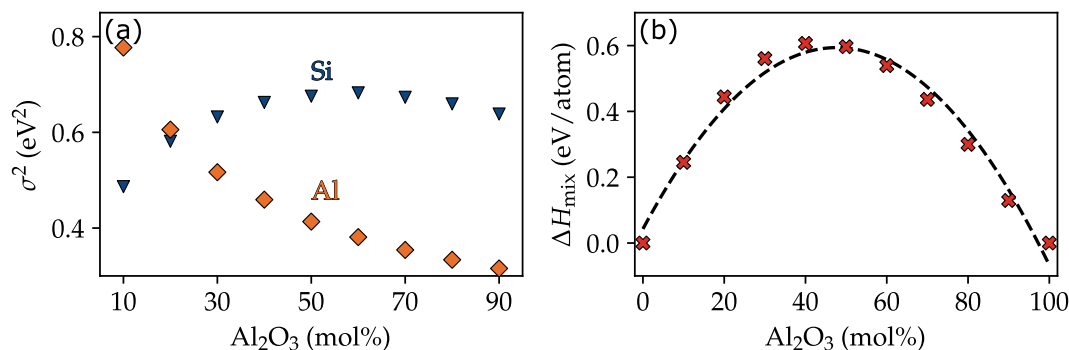


Fig. 4. (a) The variance of atomic potential energy for Si and Al atoms. (b) Variation of the mixing enthalpy as a function of alumina content at 300 K. The lines in (b) are to guide the eye (For interpretation of the colors in the figure(s), the reader is referred to the web version of this article).

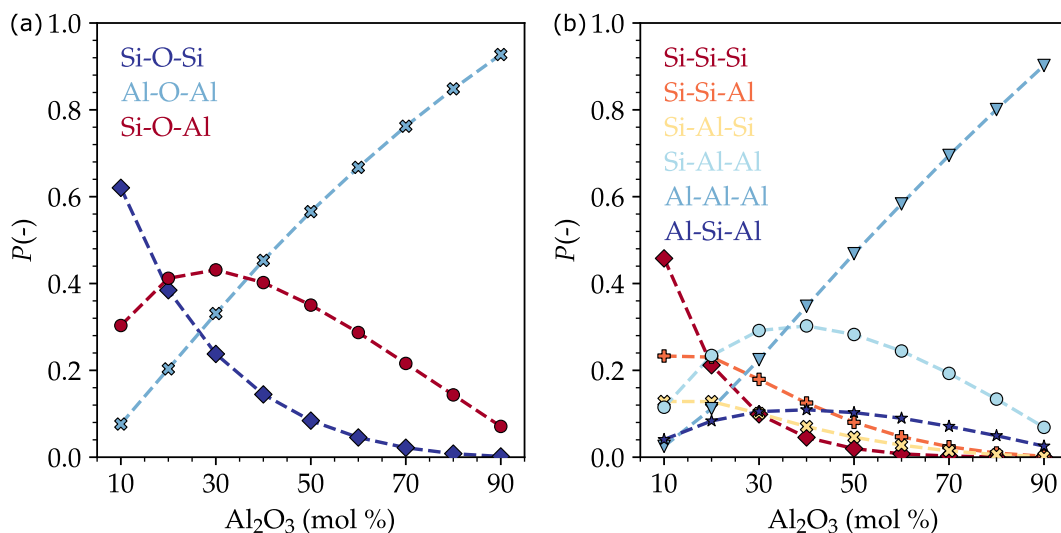


Fig. 5. (a) T-O-T and (b) T-T-T linkages (T is Si or Al) as a function of the Al<sub>2</sub>O<sub>3</sub> content in the MD simulated AS glasses at 300 K.

mol% of Al<sub>2</sub>O<sub>3</sub>, indicating the local environment of Al becomes more homogeneous with increasing Al<sub>2</sub>O<sub>3</sub> content.

On the other hand, this observation holds physical significance: the environment of SiO<sub>4</sub> tetrahedra is more homogeneous at low alumina content. With increasing Al<sub>2</sub>O<sub>3</sub>, the values of the variance increase to reach a maximum of around 50 mol% of Al<sub>2</sub>O<sub>3</sub> and then decrease again. This also indicates that the highest deviation from the mean potential energy of Si atoms occurs when the Al<sub>2</sub>O<sub>3</sub> content is approximately 50 mol%. This might be associated with the emergence of zones characterized by elevated potential energy. Such regions could signify areas where atomic interactions are less favorable, forming distinct phases. Thus, the phase separation in aluminosilicate glasses is due to a non-random distribution of the SiO<sub>4</sub> and AlO<sub>n</sub> polyhedra, caused by the increasing attraction between SiO<sub>4</sub> tetrahedra and repulsion between SiO<sub>4</sub> tetrahedra and AlO<sub>n</sub> polyhedra. These interactions between the network and the former polyhedra lead to the formation of distinct phases to minimize the system energy. This behavior is well captured by the positive mixing enthalpy  $\Delta H_{\text{mix}}$  (detail of the calculation of  $\Delta H_{\text{mix}}$  is given in the SM) obtained from the MD simulations, as shown in Fig. 4(b). For all compositions, the mixing enthalpy is positive, indicating the presence of a demixing in the glass, with its maximum being around 50 mol% of Al<sub>2</sub>O<sub>3</sub>. This is also in agreement with the observations made by Jantzen et al. [49], who suggested that the modulations below T<sub>g</sub> are caused by structural density changes.

It is also known in the literature that the aluminum avoidance principle was proven to be violated in many aluminosilicate glasses, specifically when there is not enough Si in the glass to satisfy the aluminum avoidance principle [43,50]; however, if not enough Si

atoms are available to satisfy the principle Al-O-Al bonds will form. These cases are well shown in Fig. 5(a), where we show the change of T-O-T with T is Si or Al as a function of Al<sub>2</sub>O<sub>3</sub> content. The percentage of Al-O-Al linkage increases, while that of Si-O-Si decreases with increasing Al<sub>2</sub>O<sub>3</sub> content. Moreover, in Fig. 5(b), the percentage of T-T-T linkages where T is Si or Al is depicted. The statistical analysis of T-T-T linkages demonstrates that with the increase of Al<sub>2</sub>O<sub>3</sub>, the Si-Si-Si, Si-Si-Al, and Si-Al-Si linkages decrease as a function of Al<sub>2</sub>O<sub>3</sub> and the Al-Al-Al linkages increase linearly with alumina content. On the other hand, Si-Al-Al and Al-Si-Al increase to reach a maximum at around 40 mol% and decrease again. This supports the results observed from the WC parameter, where initially the glass is made of a Si-rich phase dominated by SiO<sub>4</sub> linkages, and as the Al<sub>2</sub>O<sub>3</sub> content increases, an Al-rich and a mixed Si and Al phase coexist, which means the occurrence of phase separation. With further increase of Al<sub>2</sub>O<sub>3</sub>, the AlO<sub>n</sub> phase dominates, which leads to the formation of a single phase again. It is to be stressed that the phase separation is mainly driven by a competition between entropy and enthalpy that force atoms to have the arrangement depicted above. This competition in the simulated glasses depends on the interatomic potential and the cooling rate used, which might lead to a shift of the maximum phase separation to high Al<sub>2</sub>O<sub>3</sub> content.

In summary, the current study has contributed to a clearer understanding of the origins of nanoscale phase separation in aluminosilicate glasses. We showed that compositions containing approximately 50 mol% of alumina exhibited the most pronounced phase separation. The origins of the nanoscale phase separation are caused by a preference for

SiO<sub>4</sub> and AlO<sub>n</sub> polyhedra to cluster starting from the second neighboring shells, ultimately leading to the formation of distinct, homogeneous phases locally. This was demonstrated through the chemical order analysis at larger coordination shells using the WC parameter. Moreover, we showed that the phase separation is particularly pronounced in the vicinity of the glass transition temperature and continues to be notable at lower temperatures.

### CRedit authorship contribution statement

**Houssam Kharouji:** Writing – original draft, Methodology, Investigation, Formal analysis. **Abdellatif Hasnaoui:** Writing – review & editing, Supervision. **Achraf Atila:** Writing – review & editing, Writing – original draft, Visualization, Methodology, Investigation, Formal analysis, Validation, Data curation, Conceptualization, Supervision.

### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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### Appendix A. Supplementary data

Supplementary material related to this article can be found online at <https://doi.org/10.1016/j.mtla.2024.102148>.

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